

On the $H\alpha$ Behaviour of Blue Supergiants: Rise and Fall over the Bi-stability Jump

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ABSTRACT

Context. The evolutionary state of blue supergiants is still unknown. Stellar wind mass loss is one of the dominant processes determining the evolution of massive stars, and it may provide clues on the evolutionary properties of blue supergiants. As the $H\alpha$ line is the most oft-used mass-loss tracer in the OB-star regime, we investigate $H\alpha$ line formation as a function of T_{eff} .

Aims. We provide a detailed analysis of the $H\alpha$ line for OB supergiant models over an T_{eff} range between 30 000 and 12 500 K, with the aim of understanding the mass-loss properties of blue supergiants.

Methods. We model the $H\alpha$ line using the non-LTE code CMFGEN, in the context of the bi-stability jump at $T_{\text{eff}} \sim 22\,500$ K.

Results. We find a maximum in the $H\alpha$ equivalent width at 22 500 K - *exactly* at the location of the bi-stability jump. The $H\alpha$ line-profile behaviour is characterised by two branches of effective temperature: (i) a *hot branch* between 30 000 and 22 500 K, where $H\alpha$ emission becomes stronger with decreasing T_{eff} , and (ii) a *cool branch* between 22 500 and 12 500 K, where the $H\alpha$ line becomes *weaker*. Our models show that this non-monotonic $H\alpha$ behaviour is related to the optical depth of $\text{Ly}\alpha$, finding that at the “cool” branch the population of the 2nd level of hydrogen is enhanced in comparison to the 3rd level. This is expected to increase line absorption, leading to weaker $H\alpha$ flux when T_{eff} drops from 22 500 K downwards. We also show that for late B supergiants (at T_{eff} below $\sim 15\,000$ K), the differences in the $H\alpha$ line between homogeneous and clumpy winds becomes insignificant. Moreover, we show that at the bi-stability jump $H\alpha$ changes its character completely, from an optically thin to an optically thick line, implying that *macro-clumping* should play an important role at temperatures below the bi-stability jump. This would not only have consequences for the character of observed $H\alpha$ line profiles, but also for the reported discrepancies between theoretical and empirical mass-loss rates.

Key words. stars: mass loss – stars: supergiants – stars: atmospheres – stars: evolution – stars: winds, outflows

1. Introduction

Blue supergiants (BSG) are key creators of heavy elements, thereby contributing to the composition of the interstellar medium; they are visible out to large distances, and dominate the spectra of star-formation galaxies. Moreover, they heat the interstellar gas and dust and produce the far-infrared luminosities of galaxies. B supergiants (Bsgs) in particular are indissoluble troublemakers in today’s Astrophysics, as a proper understanding of their evolution is still in its infancy (Langer 2012). Even their fundamental properties are not yet known, and the basic issue as to whether they are core hydrogen (H) burning main-sequence (MS) or core helium (He) burning post-MS stars is still under debate (Vink et al. 2010; Georgy et al. 2013).

Lamers et al. (1995) reported the existence of a drop in terminal wind velocities of Bsgs by a factor of two, which is referred to as the “bi-stability jump” (Pauldrach & Puls 1990). According to model predictions of Vink et al. (1999) it is expected that the faster winds of hotter stars switch to slower winds with a ~ 5 times higher \dot{M} below $T_{\text{eff}} \sim 22\,000$ K. Whereas the drop in terminal velocities has been confirmed by observations (Markova & Puls 2008), the predicted increase in \dot{M} is still controversial.

Whilst Benaglia et al. (2007) and Markova & Puls (2008) uncovered a local maximum in both radio and $H\alpha$ mass-loss rates at the location of the bi-stability jump, several works have highlighted significant discrepancies between theoretical and empirical mass-loss rates for B1 and later B-type supergiants (Vink et al. 2000; Trundle et al. 2004; Trundle & Lennon 2005;

Crowther et al. 2006; Benaglia et al. 2007; Markova & Puls 2008). For the late Bsgs Vink et al. (2000) noted that their predictions agreed reasonably well with empirical rates from both radio and $H\alpha$ emission, but they found huge (order of magnitude) discrepancies when the $H\alpha$ line was PCyg shaped or in absorption. Searle et al. (2008) emphasised that empirical models for Bsgs likely have an incorrect ionisation structure as they found it challenging to reproduce the optical $H\alpha$ line simultaneously with key ultraviolet (UV) diagnostics.

To make the picture even more complex, one should be aware of discrepancies between mass-loss rates estimated from $H\alpha$, UV, and radio observations for OB stars in general (Massa et al. 2003; Bouret et al. 2005; Puls et al. 2006; Fullerton et al. 2006), which may be due to distance-dependent wind clumping and/or porosity effects (Oskinova et al. 2007; Sundqvist et al. 2010, 2011; Muijres et al. 2011; Šurlan et al. 2012).

In any case, the general trend seems to indicate $\dot{M}_{\text{vink}} > \dot{M}_{H\alpha}$ for B1 and later supergiants, whilst the reverse holds for earlier O supergiants. The key question is whether this discrepancy is the result of incorrect predictions or alternatively that we may not understand the mass-loss indicator $H\alpha$ well enough. What would one expect to happen when FeIV recombines, and FeIII starts to control the wind driving (Vink et al. 1999)?

Over the last decade we have made a dedicated effort to improve the physics in the Monte Carlo line driving calculations. We now solve the wind dynamics more locally consistently (Müller & Vink 2008; Muijres et al. 2012), we added Fe to the statistical equilibrium calculations in the ISA-WIND model at-

mosphere (de Koter *et al.* 1993) (rather than treating this important line-driving element in a modified nebular approximation), and finally, we have studied the effects of wind clumping and porosity on the line driving (Muijres *et al.* 2011). After all these improvements in the Monte Carlo line-driving physics, we have to admit that the basic problem of $\dot{M}_{\text{vink}} > \dot{M}_{\text{H}\alpha}$ for B1 and later supergiants, is still present, and it is time that we also consider the possibility that it is not the predictions that are at fault, but that we do not understand H α line formation in Bsgs sufficiently well to allow accurate mass-loss determinations from H α .

In order to obtain a more complete picture of H α as a mass-loss diagnostic, we need to know not only the clumping properties of BSGs, but also how sensitive H α is to the clumping on both sides of the bi-stability jump. Currently, T_{eff} dependence of the mass-loss rates from H α of BSGs is still uncertain, and a better knowledge is required, especially as it impacts the question of the evolutionary nature of BSGs.

Despite the fact that the spectral modelling of Bsgs is an active area of research (Zorec *et al.* 2009; Fraser *et al.* 2010; Castro *et al.* 2012; Clark *et al.* 2012; Firnstein & Przybilla 2012), due to a lack of observational data, the nature of Bsgs is a long-standing issue in stellar evolution. Hunter *et al.* (2008) argued on the basis of high-resolution VLT-FLAMES data that the slowly rotating Bsgs in the Large and Small Magellanic Clouds (LMC; SMC) are post-MS objects (although their large numbers would remain unexpected). Vink (2008) and Vink *et al.* (2010) established a general lack of fast rotating Bsgs ($v \sin i > 50 \text{ km s}^{-1}$) and suggested that these objects could naturally be explained as MS stars, if they lose their angular momentum via an increased mass-loss rate due to the bi-stability jump (the so-called bi-stability braking). They also pointed out that this mechanism would be efficient if the stars spend significant amounts of time on the MS. While the first hypothesis received some support by the apparently brightest supernova (SN) in the telescopic era: SN 1987A, bi-stability braking for stars with initial masses above $40 M_{\odot}$ was confirmed by Markova *et al.* (2014).

Hot stars with high mass-loss rates are expected to have strong emission lines, predominantly formed by recombination. The H α line is an excellent tracer for hot-star mass loss, as its detection is possible for large numbers of stars. Moreover, the strong velocity dependence of H α combined with high-resolution spectroscopy could provide valuable information about velocity fields and density structures in Luminous Blue Variables (LBVs) and SN progenitors (Groh & Vink 2011). That is why understanding the behaviour of the H α line is crucial to our understanding of mass loss.

After this introduction, in Sect. 2, we describe shortly the used input parameters for our supergiant models. In Sect. 3, we show that the computed line profiles exhibit non-monotonic temperature behaviour, and provide the most likely explanation for it. Section 4 explores the impact of clumping on H α line equivalent width (EW) and optical depth. In Sect. 5 we discuss the importance of our findings and finally in Sect. 6 we summarise our results and discuss possible implications for future work.

2. Method and input model parameters

In order to study H α line formation, we use the non-LTE radiative transfer code CMFGEN (Hillier & Miller 1998; Hillier & Lanz 2001), and we calculate a model grid over a range of temperature and $\log g$ appropriate for Bsgs (cf. Table 1). Apart from changes in T_{eff} and $\log g$, H α line formation is sensitive to luminosity, L_{\star} , mass, M_{\star} , mass-loss rate \dot{M} , clumping, and velocity structure. As CMFGEN does not solve the momentum equation of

Table 1. Adopted stellar parameters used in the model grid.

$\log\left(\frac{L_{\star}}{L_{\odot}}\right) = 5.5$, $M_{\star} = 40 M_{\odot}$, $\dot{M} = 2.33 \times 10^{-6} [M_{\odot} \text{yr}^{-1}]$				
T_{eff} [K]	R_{\star} [R_{\odot}]	$v_{\infty} = 2 \times v_{\text{esc}} [\frac{\text{km}}{\text{s}}]$	$\log(g)$	model
30 000	21	1701	3.40	M
27 500	25	1558	3.25	
25 000	30	1447	3.09	
22 500	37	1276	2.90	
20 000	46	1138	2.70	
17 500	61	998	2.47	
15 000	83	855	2.20	C
12 500	120	715	1.88	

the wind, a velocity structure needs to be assumed. For the accelerating part of the wind, we adopt a standard β -type velocity law with $\beta = 1$, whilst a hydrostatic solution is adopted for the subsonic region.

To first understand how the H α line profile changes in a *qualitative* sense, we use models with a simplistic H chemical composition, keeping L_{\star} , M_{\star} , $v_{\infty}/v_{\text{esc}}$ and \dot{M} fixed, whilst we vary T_{eff} over a range from 30 000 K down to 12 500 K. Following Vink *et al.* (2001), we adopt the parameters listed in Table 1.

Note that for fixed L_{\star} , the changes in T_{eff} inevitably lead to different R_{\star} and v_{esc} values, which may influence the H α EW (to account for this, we used an optical depth parameter tool that was introduced by Puls *et al.* (1996); see Sect. 3.1). To keep as many model parameters as possible fixed, we used $v_{\infty}/v_{\text{esc}} = 2$ – the mean value of $v_{\infty}/v_{\text{esc}}$ ratio at both sides of the bi-stability jump (Lamers *et al.* 1995; Markova & Puls 2008). It is more natural to keep the $v_{\infty}/v_{\text{esc}}$ ratio constant rather than the value of v_{∞} itself, as the models have different radii.

It is important to keep in mind that due to metal line blocking and a blend with He II the H α line could be sensitive to changes in He as well as metal abundance. To estimate these effects we also compare how the H α EW behaves in models with different chemical complexities: (i) pure H, (ii) H + He, with two different He mass fractions ($Y_{\text{He}} = 0.25$ and 0.6), as well as (iii) metal line-blanketed models.

3. Results

3.1. H α line profile and equivalent width

Our systematic examination of the H α line for supergiant models over T_{eff} range between 30 000 and 12 500 K shows non-monotonic changes in the H α line profiles with T_{eff} . As shown in Fig. 1, the H α line becomes stronger when T_{eff} drops from 30 000 to 22 500 K, where the line reaches its peak value. Below 22 500 K the line becomes weaker when the effective temperature is further reduced to 12 500 K.

Although we have kept all stellar parameters fixed, the models have different radii and v_{∞} . Note that in the hottest models the radii are up to a factor of 6 smaller, whilst the terminal velocities are only up to a factor 2.4 higher. To extract the *true* temperature effect, the different radii and terminal velocities should be taken into account. One way to do this is through the use of the wind-strength parameter Q concept introduced by Puls *et al.* (1996). It was demonstrated that for O-type stars recombination lines remain unchanged by specific sets of individual values of \dot{M} , v_{∞} and R_{\star} as long as the wind-strength parameter Q :

$$Q = \frac{\dot{M}}{(v_{\infty} R_{\star})^{1.5}} \quad (1)$$

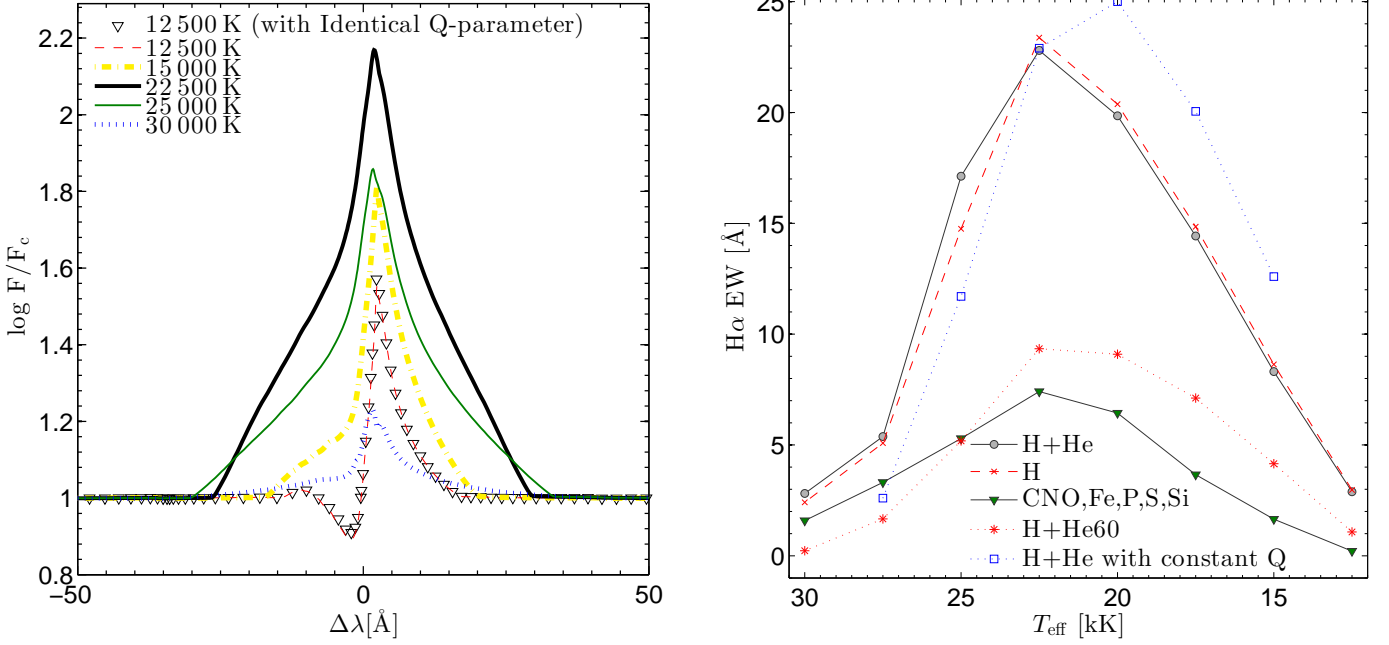


Fig. 1. *Left:* H α line profiles for CMFGEN models with parameters as listed in Table 1. Black triangles represent the line profile with the same Q -parameter (Eq. 1) as model C ($T_{\text{eff}}=12\,500$ K) but with different \dot{M} and R_* values. *Right:* H α line EW vs T_{eff} for models with only H (crosses), H+He (circles) and more sophisticated (triangles) models. Red asterisks represent the changes in the H α line when the He mass fraction in the pure H+He models is increased to 60%. Blue squares indicate how the H α EW behaves as a function of a constant Q value.

is invariant and T_{eff} remains unchanged. Here \dot{M} is calculated in M_\odot/yr , v_∞ in km s^{-1} and R_* in R_\odot . As an example, in Fig. 1 we show that the H α line profile from model C (red dashed line) is basically unaffected when the mass-loss rate is decreased by factor of two and the radius by $2^{2/3}$ (black triangles).

The right-hand panel of Fig. 1 displays how the H α EW behaves as a function of T_{eff} for models with similar parameters, but with different individual \dot{M} which scale in such a way that Q is constant (blue squares). The peak is still present and this implies that the H α behaviour is not due to the changes in $v_\infty \times R_*$, i.e., it is a real temperature effect. Note that the peak is slightly shifted towards $T_{\text{eff}} \sim 20\,000$ K.

To gain additional insight, the EW was separated into the line and continuum flux (respectively illustrated in Fig. 2 with red circles and squares). The reason for doing so is that the EW is a measurement for the ratio of the line flux over the continuum flux, and models with similar EW may have completely different line fluxes. As can be seen in Fig. 2, the coolest model has a seven times larger line flux than the line flux in the hottest model, whilst in Fig. 1 it seems that both models have a similar line strength. The similar H α EW in both models is due to a constantly increasing continuum flux when reducing T_{eff} , and the higher line flux in the coolest model. The peak in right-hand panel of Fig. 1 results at 22,500 K (and not at 20,000 K as in Fig. 2) because the ratio between the line and continuum fluxes is the largest for that T_{eff} . Note that the cooler supergiants models have larger radii which increases the continuum flux with decreasing T_{eff} .

The right-hand panel of Fig. 1 also presents behaviour patterns in models with different chemical complexities. The simplest set of models (with only H; i.e. without He) presented with

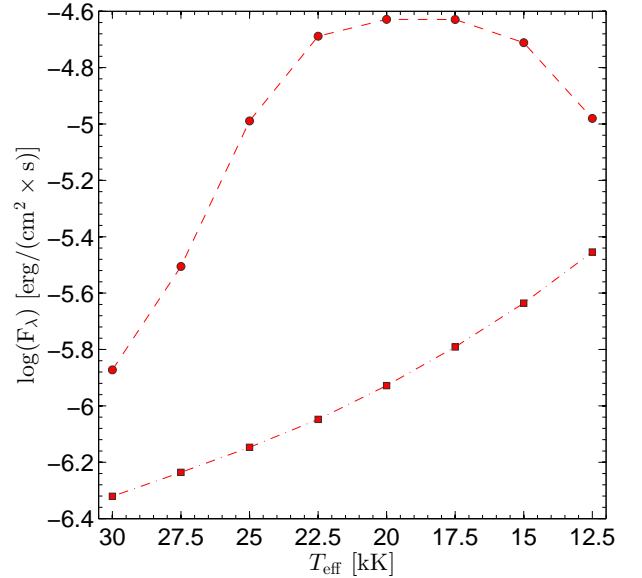


Fig. 2. Integrated line (circles) and continuum flux (squares) at the wavelength of H α for H + He models. Note that the flux represents the flux at a distance of 10 parsec.

red crosses show practically identical H α EWs¹ as the EWs from models including both H and He (displayed with grey circles). In other words, He does not seem to influence the H α line. The reason could be that for this T_{eff} range HeII is diminished and the blueward HeII blend is not essential, whilst the HeI continuum plays only a minor role at these temperatures. However, the lat-

¹ We have defined the line EW to be positive for an emission line and negative for an absorption line.

Table 2. Atomic data used for our simplistic H + He supergiant models. For each ion the number of full levels, super levels, and bound-bound transitions are provided.

Ion	Super levels	Full levels	b-b transitions
H I	20	30	435
He I	45	69	905
He II	22	30	435

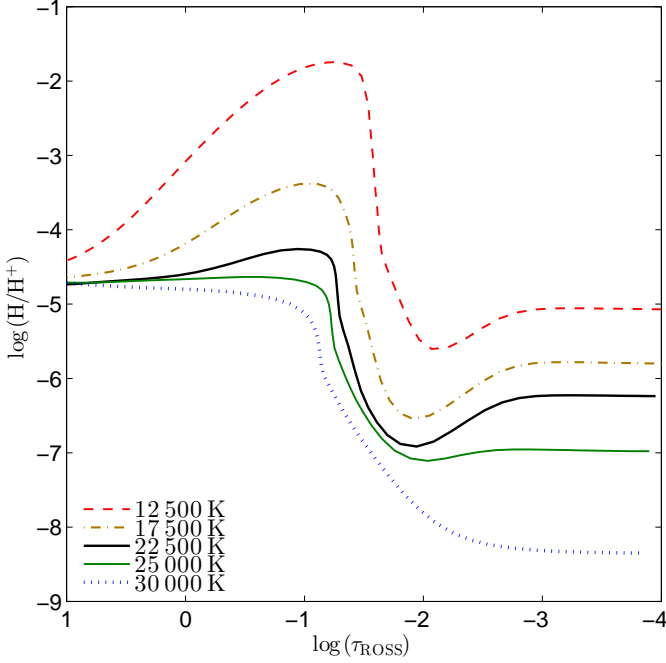


Fig. 3. Hydrogen ionisation structure for models with various T_{eff} .

ter holds only when the He abundance is less than 25% by mass. The red asterisks in the right-hand panel demonstrate that the He rich models (comprising 60% He mass fraction) have a factor 2-3 lower $H\alpha$ EW. This behaviour is in agreement with the results of Dimitrov (1987) who found that a high abundance of He produces stronger absorptions in H lines. Despite the quantitative differences, He is not important for the *qualitative* behaviour of $H\alpha$ versus T_{eff} .

The line EW values from the simplistic H+He models (grey circles) are also compared to those determined from more realistic chemical models (green triangles), which include C, N, O, Si, S, P and Fe (with atomic data as listed in the table from Appendix B and half-solar metal abundances). Whilst there are quantitative differences due to line blanketing, the *qualitative* behaviour in $H\alpha$ EW is similar in both sets of models with different complexities. This implies that the reason for their behaviour is fundamentally driven by the properties of H. Despite their simplification, the models including H+He only (with atomic data as indicated in Table 2), provide an overall picture of the effective temperature dependence of $H\alpha$ for Bsgs. It is therefore reasonable to take advantage of these H+He models, using them as a starting point for our investigation.

As an aside, we found that for models with a flatter velocity law ($\beta = 2$) or with different mass-loss rates (varying from $\sim 10^{-7}$ to $10^{-5} M_{\odot} \text{yr}^{-1}$) the resulting $H\alpha$ EW behaviour was qualitatively similar to those presented in Fig. 1 with $H\alpha$ EW showing a peak at $T_{\text{eff}} \approx 22\,500$ K.

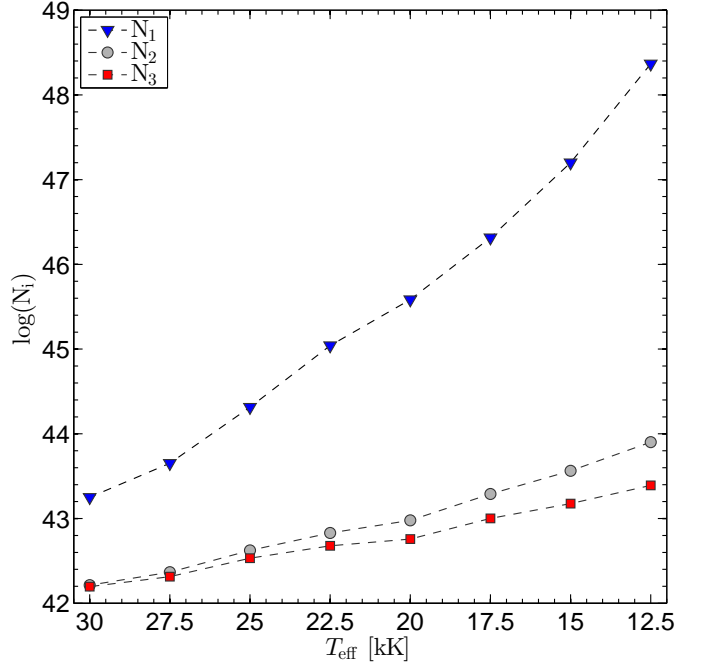


Fig. 4. Total number of H atoms in the stellar wind versus T_{eff} . Note that the total number of H atoms is determined from $\tau_{\text{ross}} < 2/3$.

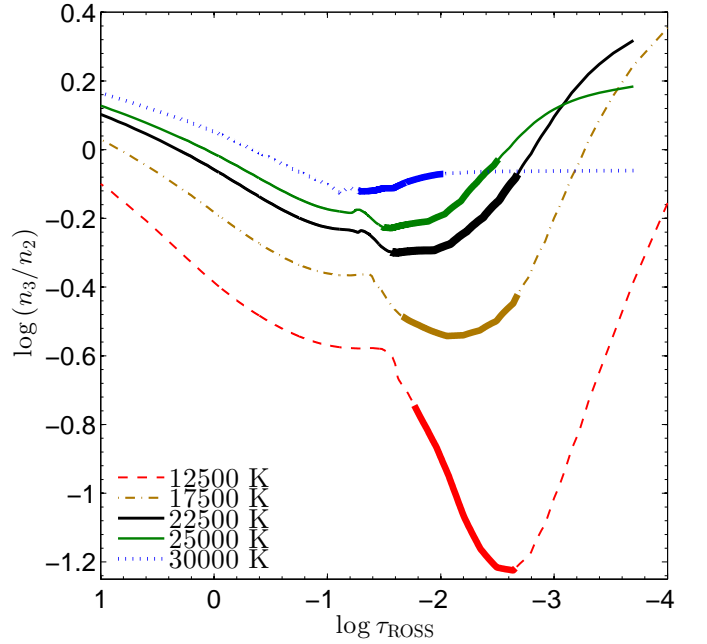


Fig. 5. Changes in the (n_3/n_2) ratio with T_{eff} . Regions where most of the emergent $H\alpha$ photons originate from are represented with a thick solid line (cf. Appendix A).

3.2. Two branches of $H\alpha$ behaviour

The $H\alpha$ EW changes drastically with T_{eff} in all sets of models. In general, we find there to be two branches in Figs. 1 and 2: the “hot” and “cool” ones. The hot branch is located between 30 000 to 22 500 K, where the $H\alpha$ line emission becomes stronger with decreasing T_{eff} . At the cool branch, from 22 500 to 12 500 K, the behaviour of the line flux changes in the opposite direction. This

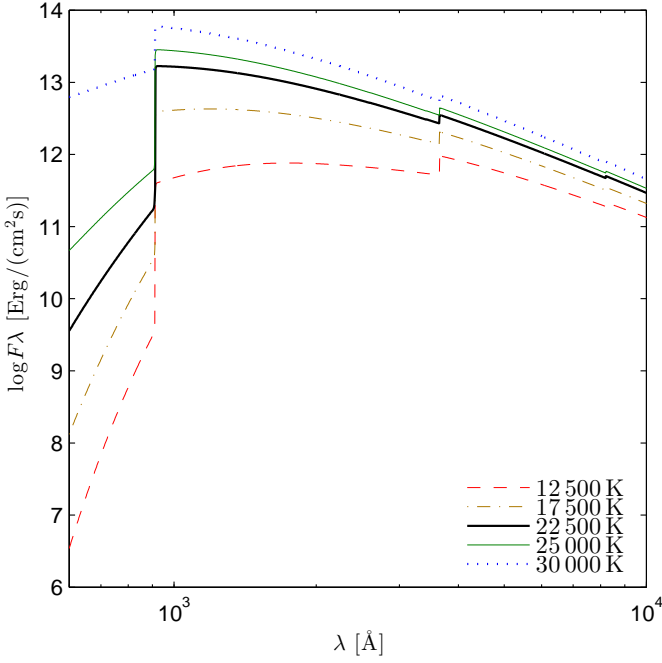


Fig. 6. Spectral energy distribution at the stellar surface of our models.

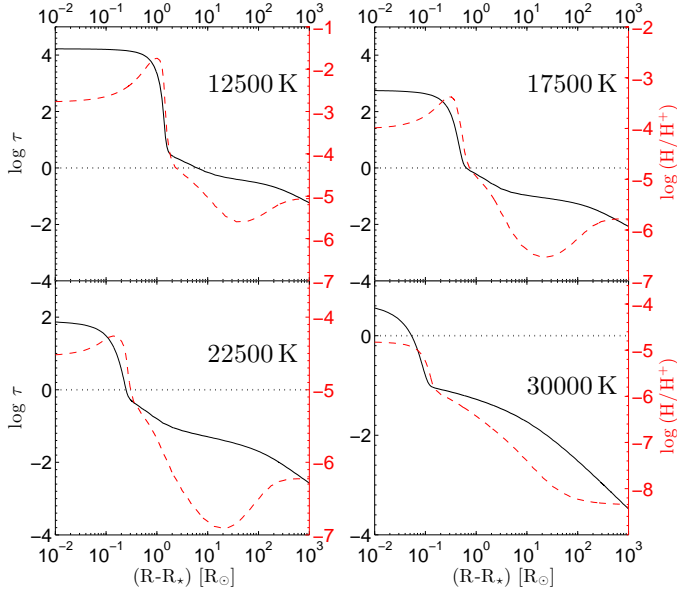


Fig. 7. Comparison between the H ionisation structure (red dashed line) and the Lyman continuum optical depth at $\lambda \sim 900$ Å (black solid line) versus the distance from the stellar photosphere. Solid lines are reserved for the wind optical depth, whilst the dotted horizontal lines indicate the transition between optically thick and thin part of the wind in the Lyman continuum ($\tau = 1$). Red colour on the right-hand side is used for the H ionisation structure.

implies that there is a qualitative change in the behaviour of H α around 22 500 K.

3.2.1. The “hot” branch

The formation of the H α line involves transitions between the 3rd and 2nd level of H. Therefore, the emission should be propor-

tional to the total number of H atoms in the 3rd level in the wind above $\tau_{\text{ROSS}} = 2/3$ (N_3). As N_3 (as do the number of H atoms in other levels) scales with the fraction of neutral H, we show the H ionisation structure of our models in Fig. 3. It is evident that the wind is mostly ionised in all models and as T_{eff} drops the wind recombines slightly. The fraction of neutral H increases to almost two percent when T_{eff} is reduced to 12 500 K (at $\log \tau_{\text{ROSS}} \sim -1.5$, $\log (H/H^+) \sim -1.8$). Furthermore, in an absolute sense the total number of H atoms in the second level (N_2) and N_3 increase, as illustrated in Fig. 4. As a result, the flux in H α , which is proportional to N_3 , should to first order increase as T_{eff} drops. Therefore, the trend in H α on the “hot” side of Figs. 1 and 2 can be understood in terms of a simple recombination effect (see also Puls *et al.* 1996). Difficulties with such a simple explanation would arise if we were to try to explain the existence of the cool branch in a similar way. Figure 4 illustrates that when T_{eff} is reduced from 22 500 to 12 500 K, N_3 is still increasing and this should probably produce a stronger H α line for cooler models. Contrary to expectation, the opposite behaviour of H α is produced for this branch.

3.2.2. The “cool” branch

On the cool branch, small changes in T_{eff} lead to qualitatively different H α line profiles. Most notable is the appearance of a P Cygni profile in Fig. 1 when the effective temperature is reduced from 15 000 to 12 500 K. In fact, these differences between both H α line profiles contain major clues to the unexpected changes in the line flux of the models on the cool branch. In case H α were a pure recombination line only the third level would be relevant. However, as shown later, H α increases its optical depth at cooler T_{eff} . It is thus necessary to assess the source function and the ratio of the number of H atoms per cm³ in the 3rd over the 2nd level (n_3/n_2).

Figure 4 reveals that N_3 is constantly increasing when T_{eff} is reduced, however, N_2 increases more steeply. The (n_3/n_2) ratio has been plotted versus Rosseland optical depth in Fig. 5. First of all, the (n_3/n_2) ratio is close to unity for the models at the “hot” branch. However, for the models on the cool branch, n_2 becomes significantly larger than n_3 , particularly in the outer wind. Although N_1, N_2, N_3 are all increasing as T_{eff} drops, the *ratio* (n_3/n_2) is always decreasing, and the second level becomes more populated than the third level on the cool branch. This leads to a “dip” in (n_3/n_2) in the outer wind of the cooler models. The “dip” is quite pronounced, especially for the coolest model, where n_2 is ten times higher than n_3 . The increased n_2 in the outer wind produces absorptions for H α photons emitted close to the photosphere. This naturally decreases the H α flux. Hence the question about the decreasing EW over the cool branch, could be referred to as an issue regarding the behaviour of the (n_3/n_2) *ratio*.

Therefore, the next question is why n_2 becomes significantly larger than n_3 ? Is the “dip” in (n_3/n_2) ratio is predominantly caused by an *increase* of n_2 or by a *decrease* of n_3 ? In order to address these questions, we take a detailed look into the behaviour of the continua and their effects on individual levels.

3.3. Two possible explanations for the existence of the “cool” branch

3.3.1. A decrease of n_3 ?

In Fig. 6 we show spectral energy distributions (SEDs) at the surface. The Lyman continuum flux is greatly reduced from 30 000

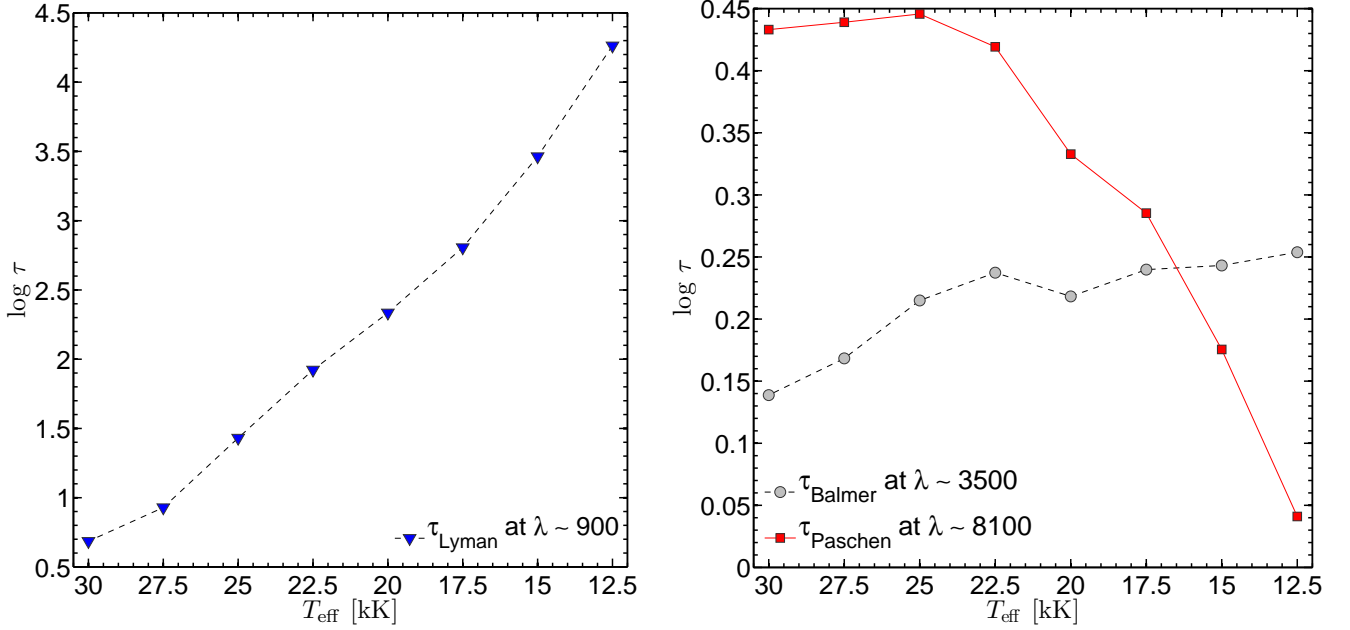


Fig. 8. Wind optical depth at $\tau_{\text{ross}} = 2/3$ in the Lyman (left), Balmer (grey circles) and Paschen continua (red squares) (right).

to 12 500 K. Despite that, nothing dramatically happens with the Lyman continuum around 22 500 K - the changes in the continuum flux are rather gradual. To understand the behaviour of the Lyman continuum in Fig. 7 we compare the continuum optical depth at $\lambda \sim 900$ Å (black solid line) to changes in the H ionisation structure (red dashed line with ordinate in red colour placed on the right side). It is evident that the Lyman continuum becomes optically thick at distances closer than $d_r \sim 0.06, 0.25, 0.6$ and $7 [R_\odot]$ from the photosphere (roughly $\tau_{\text{ross}} = 2/3$), respectively for the models with $T_{\text{eff}} = 30\,000, 22\,500, 17\,500$ and $12\,500$ K. The fraction of neutral H at those distances is between $\sim 10^{-4.9}$ and $10^{-4.3}$. The comparison between the Lyman continuum optical depth and H ionisation structure shows that as soon as neutral H atoms exceed *critical* values, the Lyman continuum becomes optically thick. Moreover, the steep increase of neutral H close to the star leads to a significant increase of the optical depth of the Lyman continuum. Consequently, a large fraction of Lyman ionising photons are blocked, and the Lyman continuum is no longer the main source of ionisation.

Najarro *et al.* (1997) studied the appearance of H and He lines in the wind of the LBV P Cygni, and they found that H recombination, crucial for the Ly continuum and the Ly α optical depth, block the ionising Lyman flux. In their models, the increased n_2 is due to the high optical depth of Ly α , giving rise to strong absorption, similar to those produced in our coolest B supergiant model.

As n_1, n_2 and n_3 are controlled from different continuum ranges, knowledge about the behaviour of Balmer and Paschen fluxes (not only of the Lyman continuum) is required to understand their behaviour. In Fig. 6 it is shown that the Balmer and Paschen continuum fluxes decrease when T_{eff} is reduced. To understand this, we plot the wind optical depth in the Balmer and Paschen continua in Fig. 8 (right panel). Since the cross-section for the photo-ionisation of a H atom in quantum state n by a photon of wavelength λ is:

$$\sigma_{\text{bf}} = 1.31 \times 10^{-19} \frac{1}{n^5} \left(\frac{\lambda}{5000 \text{ Å}} \right)^3, \quad (2)$$

the opacity would be proportional to λ^3 . Therefore we have chosen to plot the wind optical depth in the Balmer and Paschen continua at wavelengths close to their corresponding jumps (respectively $\lambda \sim 3500$ and 8100 Å).

Figure 8 shows that the wind optical depth in the Balmer continuum remains fairly constant for all models. Note that when T_{eff} is reduced below 22 500 K, the wind optical depth in the Paschen continuum is steeply decreased. Models which are cooler than 17 500 K even have $\tau_{\text{Paschen}} < \tau_{\text{Balmer}}$. This should produce larger changes in the Paschen flux than in the Balmer flux. To quantify the reaction of the continuum fluxes Fig. 9 shows the number of photons - N_λ , versus T_{eff} for the Lyman, Balmer, and Paschen continua respectively. The number of photons is given by:

$$N_\lambda = 4\pi R_\star^2 \int_0^{\lambda_c} \frac{\pi F_\lambda \lambda}{hc} d\lambda \quad (3)$$

where λ_c is the wavelength boundary of corresponding continuum series; R_\star is the stellar radius. The computed numbers of photons in the Lyman and Balmer continua are in reasonable agreement (by a factor of ~ 4) with those from previous studies (Thompson 1984; Diaz-Miller *et al.* 1998) if the different radii are taken into account. Balmer and Paschen fluxes are able to ionise H atoms respectively from level 2 and 3 and regulate those levels. It is evident from Fig. 9 (left panel) that these fluxes are about 4 orders of magnitude larger than the Lyman flux at the cool branch. This is a consequence of high wind optical depth in the Lyman continuum (reported in Fig. 8) and the optically thinner wind in the Balmer and Paschen continua. Hence, the Balmer and Paschen fluxes are the main sources of ionisation of H atoms in second and third levels over the cool branch. Figure 9 indicates that the total number of photons able to ionise atoms in level 2 is nearly the same for all models. This would provide nearly the same number of H atoms ionised from the second level over both branches. Therefore, we do not expect dramatic changes in n_2 due to the Balmer continuum. By contrast, the “Paschen photons” are gradually increasing in number as T_{eff} is reduced. This is seen in the right panel of Fig. 9, where

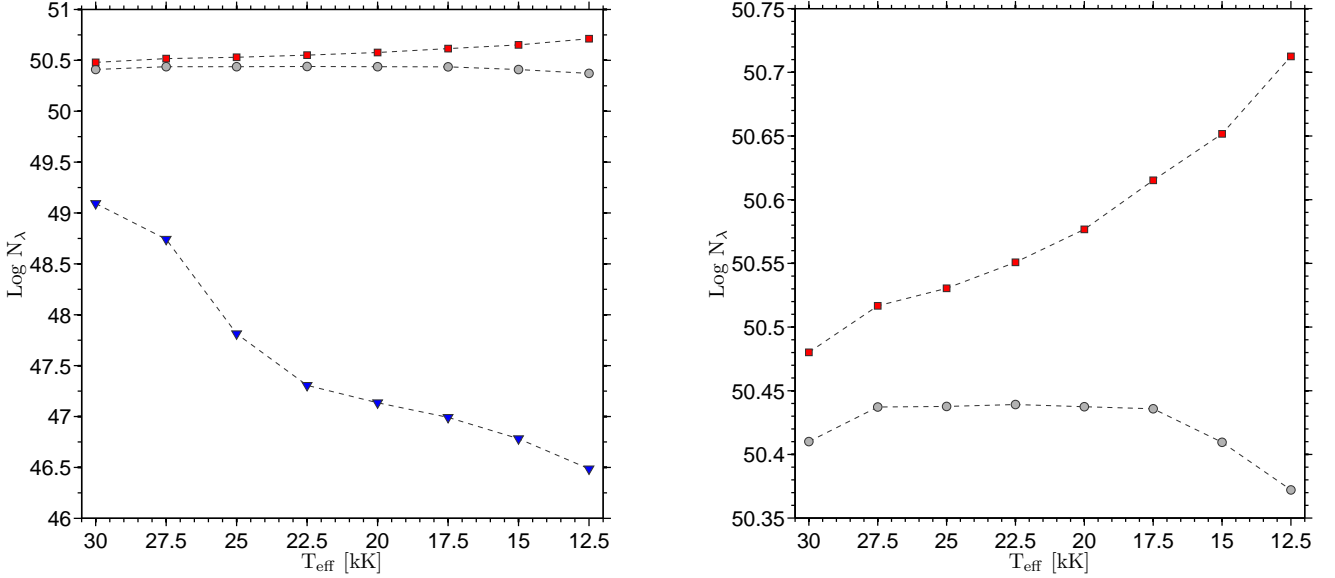


Fig. 9. Number of photons in the Lyman (blue triangles), Balmer (grey circles) and Paschen (red squares) continua vs T_{eff} . Right hand side is a “zoom in” from the left hand side.

we zoomed in around N_{λ} for Balmer and Paschen continua. For cooler models, the increasing flux in the Paschen continuum may thus depopulate more H atoms in the third level.

To understand the key question whether the “dip” in (n_3/n_2) results solely from an increase of level 2, or from a decrease of level 3 as well, we show in Fig. 10 how their number densities (and the number density of levels 4 and 10, n_4 and n_{10}) change with T_{eff} . It is evident from the plot that when T_{eff} is reduced, n_3 behaves in a fashion more similar to the number densities of the higher levels, e.g. n_4 . This is expected if the levels are mainly fed by recombinations. Therefore, a decrease of level 3 does not occur and consequently the (n_3/n_2) ratio is not affected by changes in n_3 . It seems that the changes in the Paschen continuum with T_{eff} are not large enough to cause a decrease of n_3 . As a result, n_2 should play a major role in the (n_3/n_2) ratio, causing the “dip”.

3.3.2. An increase of n_2 ?

Due to the enormous optical depth in the Ly continuum (reported in Fig. 8), the ionising flux is blocked for the “cool” models and n_1 can no longer be affected by the Lyman continuum. As a result, level 2 becomes the effective ground state, exhibiting a depopulation close to the photosphere, and an overpopulation in the outer wind in model C ($T_{\text{eff}}=12\,500$ K) as shown in Fig. 11. This is similar to the findings of Puls *et al.* (1998) who concluded that H α appears like a P-Cygni line for A supergiants. It should be noted that the produced peak in the departures from LTE for the second level (b_2) in model C, occupies the same position as the “dip” in the (n_3/n_2) ratio. This coincidence indicates that the “dip” is mainly produced by an increase of n_2 .

The difference in behaviour of level 2 compared to higher levels can clearly be seen in Fig. 10, where the number densities of level 3 (blue triangles), 4 (squares) and 10 (green solid line) are similar for all T_{eff} . For the models on the “hot” branch even n_2 behaves in an identical manner to the number densities of the higher levels. This behaviour is to be expected when the level populations are solely fed by recombinations, and when they should scale as $\rho^2 \sim r^{-4}$ (via the continuity equation). The

population of the ground state is, however, affected by ionisations as well, and it is thus inversely proportional to the dilution factor of the radiation field. The latter effect makes it increase as distance-squared, and in the outer wind the final dependence is $1/r^2$, as shown in Fig. 10. The black asterisks, connected by a dashed line, represent the number density of the ground state (n_1), which is decreasing as $1/r^2$ (illustrated as a blue dashed line). When T_{eff} is reduced to 12 500 K, n_2 diverges from higher-lying energy levels and behaves similar to the ground level (decreasing as $1/r^2$). At the same time the H α line profile forms an absorption component.

3.4. Ly α and the second level

In order to understand why level 2 diverges from the higher levels when T_{eff} drops, we ran additional models in which we artificially removed the Ly α transition by reducing its oscillator strength by a huge factor (10^4). The reaction of H α is displayed in the upper panels of Figure 12. It highlights the key role of Ly α . The effects on model M ($T_{\text{eff}}=22\,500$ K) are striking: H α now switches from a pure emission line into a PCygni line, and the line flux is decreased significantly. Furthermore, for the cooler model the removal of Ly α leads to a deeper absorption component. In the middle panels in Fig. 12 we present produced changes in n_2 (dashed) and n_3 (solid) due to the absence of Ly α – as a function of τ_{ross} . The middle panels demonstrate that the Ly α removal leads to a tremendous increase of n_2 (in comparison to the initial model) in the outer wind, leading to a stronger absorption component in model C, and the appearance of a PCygni profile in model M, where it is noted that the lack of Ly α leads to significant changes in the third level as well: n_3 is surprisingly reduced, and as a result the H α flux decreases.

To understand the changes in the level populations due the removal of Ly α , we show its net radiative rate in the lower panels of Fig. 12. This quantity is defined as the difference between the number of radiative transitions from the upper (second) level to the lower (ground) state and the number of radiative transitions from the ground level to the second level. As a result, the

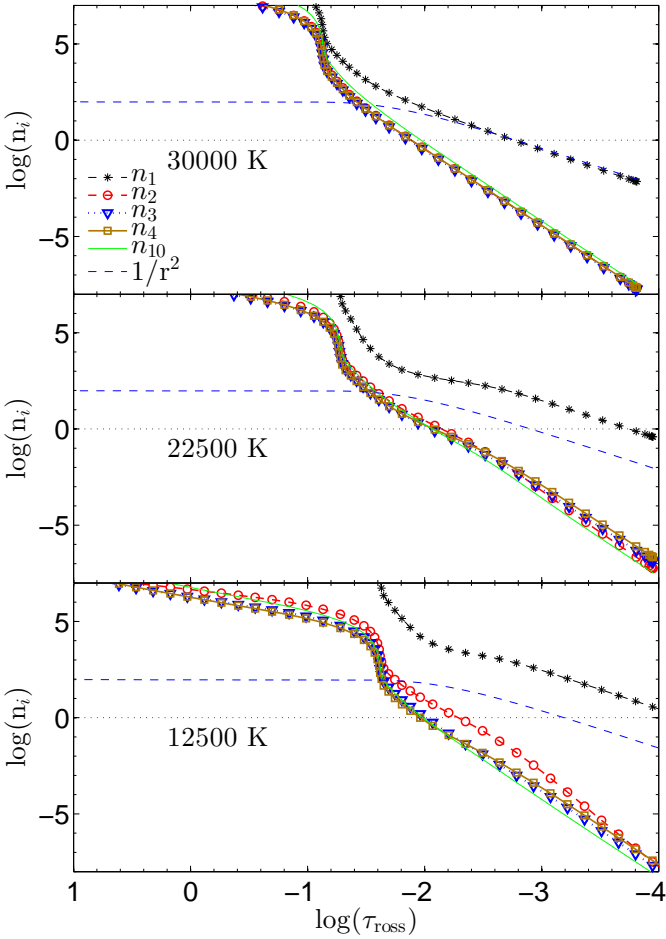


Fig. 10. Population levels of H as a function of Rosseland optical depth.

net radiative rate is positive when there is a net decay of electrons from the upper level, and it is negative when there is a net excitation of electrons. It is seen that as soon as $n_{2\text{noLy}\alpha}$ and $n_{3\text{noLy}\alpha}$ diverge from their initial ones (at $\log(\tau_{\text{ross}}) < \sim -1.8$ (in model C) and $\log(\tau_{\text{ross}}) < \sim -1.5$ (in model M)) the total Ly α radiative rate in the initial model (black solid line) is positive². Therefore, the line effectively acts as a “drain” for the second level. When we remove it, the decay of electrons from the second level is suppressed, and n_2 is tremendously increased, as shown in the middle panels.

In other words, by artificially removing Ly α , we can simulate the appearance of P Cygni profiles for hotter models, showing that the Ly α line is key to the H α behaviour. We also note that neither He nor Fe are directly required for achieving this, i.e. it is a pure H effect (as was shown in Figure 1).

Figure 13 illustrates how the Ly α optical depth changes. The location where most of the H α photons originates from (Appendix A) is shown with thicker line sections. It is evident that at the hot branch Ly α is optically thick in the inner H α forming region and becomes optically thin in the outer region. Furthermore, the Ly α optical depth at the start of the H α line-formation region is similar for hot models. However, when T_{eff} drops below 22 500 K, the Ly α optical depth at the start of the H α line-formation region steeply increases, which contin-

² Note that at $\log(\tau_{\text{ross}}) < \sim -1.8$ and $\log(\tau_{\text{ross}}) < \sim -1.5$ the H α line starts to form in models C and M (see Appendix A for details about the H α forming regions presented in Fig. 5).

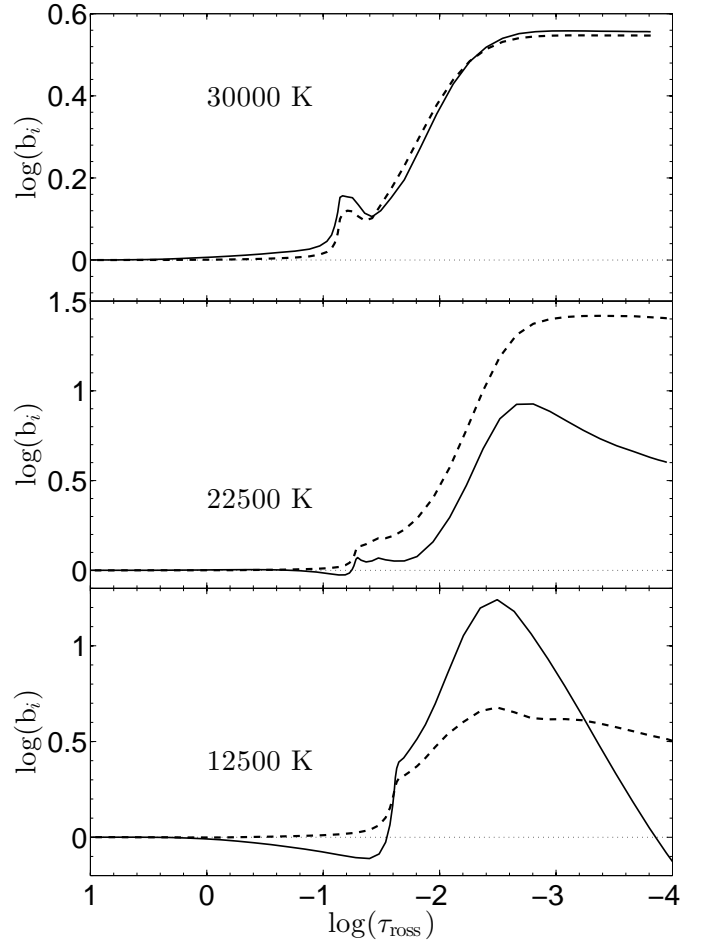


Fig. 11. Non-LTE departure coefficients for the 2nd (solid) and 3rd (dashed) level of H.

ues over the cool branch where Ly α is always optically thick throughout the entire H α line-forming region. This means that most photons from the photosphere at the Ly α wavelength do not manage to escape. I.e., they are being scattered or absorbed. As a result, the decay from the second level is effectively shut off. This can be seen in the lower panels of Fig. 12. In the region where n_2 and n_3 are affected by Ly α (i.e. for $\log(\tau_{\text{ross}}) < \sim -1.8$ for the cool model and $\log(\tau_{\text{ross}}) < \sim -1.5$ for the hotter model), the net radiative rate is 2 orders of magnitude lower for the cooler model, i.e., the second level is depopulated less efficiently by 2 \rightarrow 1 transitions.

In short, the departure of n_2 from the higher level occupation numbers (when T_{eff} drops from 22 500 K downwards) is related to the Ly α optical depth. When T_{eff} is reduced below 22 500 K, the Ly α optical depth increases steeply, and level 2 is less efficiently depleted through decay. Unlike the second level, level 3 behaves like higher levels, thus the “dip” in (n_3/n_2) is the result of an increased level 2 population.

4. The effect of clumping

Currently, CMFGEN takes optically thin (micro) clumping into account, i.e., the clumps are assumed to have a dimension smaller than the photon mean free path. The density ρ within clumps is assumed to be enhanced by a clumping factor D compared to the wind mean density $\bar{\rho}$. This factor can also be understood in terms of volume filling factor $f = D^{-1}$, assuming that the inter-

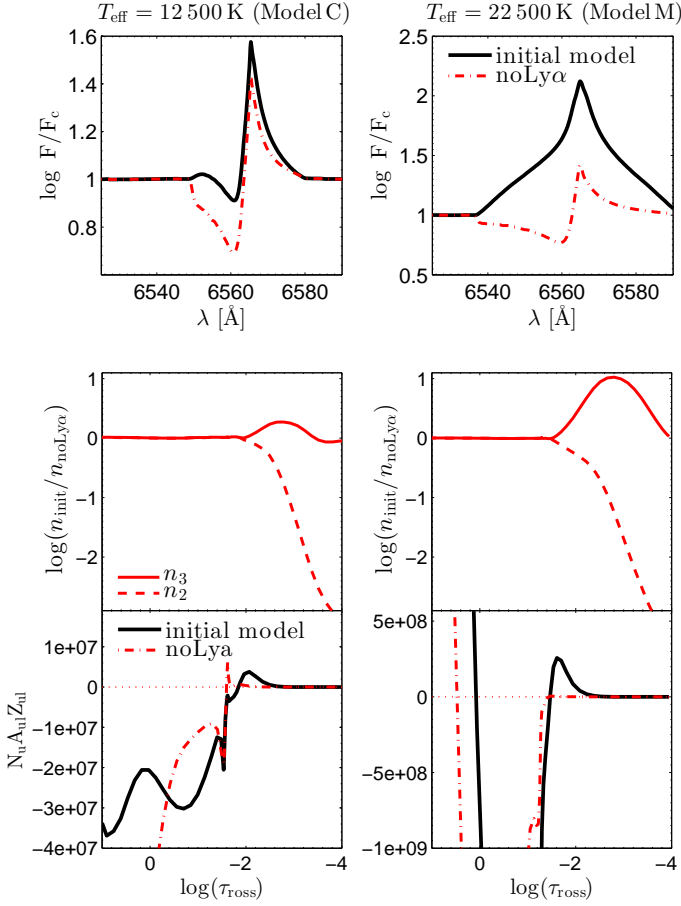


Fig. 12. *Upper panels:* effect of $\text{Ly}\alpha$ on the formation of the $H\alpha$ line: initial $H\alpha$ profile (black solid) and the profile from the models in which $\text{Ly}\alpha$ transitions were artificially removed (the red dash-dotted line). *Middle panels:* changes in the 2nd (dashed) and 3rd (solid) level of H due to the removal of $\text{Ly}\alpha$ in model C (left) and M (right). The plots present the ratio of the populations produced from the initial model over the populations from the models without $\text{Ly}\alpha$ transition. *Lower panels:* comparison of the net radiative rate of 2→1 transitions in the initial (solid black) and the model without $\text{Ly}\alpha$ (the red dash-dotted line) (see Sect. 3.4 for details).

clump medium is void. Mass-loss diagnostic techniques that are linearly dependent on density are insensitive to clumping, whilst recombination lines (sensitive to ρ^2) tend to overestimate the mass-loss rate of a clumped wind by a factor of $\sqrt{f^{-1}}$.

However, if the clumps are optically thick (macro-clumping), photons may leak through gaps between clumps, and the mean opacity is affected by the distribution, the geometrical size, and the shape of the clumps. In this case the micro-clumping approach is no longer justified.

4.1. The $H\alpha$ line in a micro-clumping approach

The $H\alpha$ line emission is a ρ^2 dependent process and is therefore sensitive to micro-clumping. In order to investigate the potential role of micro-clumping in our model sequences we calculated additional models, identical to the simplistic H+He models but with $f_\infty = 0.1$, described by the following exponential law:

$$f(r) = f_\infty + (1 - f_\infty)\exp(-v(r)/v_{\text{cl}}), \quad (4)$$

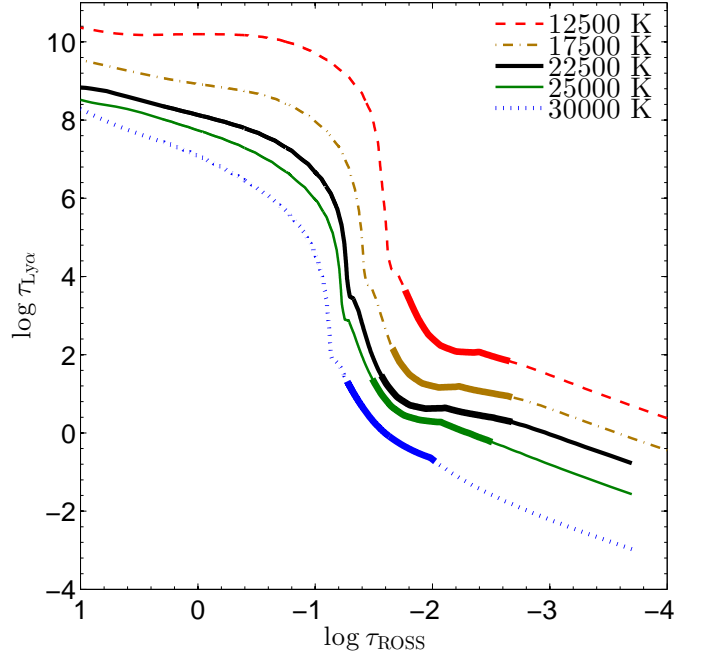


Fig. 13. Lyman α Sobolev optical depth as a function of τ_{ross} . The region where most of the emergent $H\alpha$ photons originate from are shown with thick solid lines (cf. Appendix A)

where v_{cl} is the velocity at which clumping starts. We have chosen the clumping to start at $v_{\text{cl}} = 20 \text{ km s}^{-1}$, just above the sonic point.

The $H\alpha$ line profiles are presented in the left panel of Fig. 14, where it is shown that clumping enables the asymmetry in $H\alpha$ to appear at hotter T_{eff} . When clumping was neglected, the increased level 2 produced asymmetric line profile at 15 000 K and 12 500 K (cf. Fig. 1), whilst in clumped models with $f_\infty = 0.1$, the asymmetry is already present at 20 000 K. The reason for this shift towards higher temperatures is caused by the increased mean density of the micro-clumped winds (in comparison to the mean density of smooth winds), which in turn increases the $\text{Ly}\alpha$ optical depth. Consequently, n_2 increases at hotter temperatures. The next question is whether micro-clumping may have an effect in terms of the cool versus hot branch sequences?

Therefore in the right panel of Fig. 14 we compare how the $H\alpha$ line EW behaves as a function of T_{eff} for both clumped and unclumped models. It is found that micro-clumping changes the EW dramatically in the hotter models. Also at the bi-stability jump location ($\sim 22\,500 \text{ K}$) micro-clumping has a dramatic impact on the $H\alpha$ EW. However, micro-clumping progressively plays a lesser role towards the cooler edge of the Bsg regime. The reason for this is that the second level now behaves as a quasi-ground state, i.e., it scales linearly with density ρ , and remains rather unaffected, whilst the more drastic effects for hotter models are the result of the ρ^2 scaling. Nevertheless, clumping transforms $H\alpha$ from a pure emission line into a P-Cygni line at 15 000 K. Although micro-clumping is of quantitative relevance in Bsgs (especially around the bi-stability jump), the existence of an $H\alpha$ EW peak remains present in clumped model sequences.

4.2. The $H\alpha$ optical depth in a micro-clumping approach

In Fig. 15 we compare how the $H\alpha$ optical depth changes with T_{eff} for homogeneous (left) and clumped (right) models. The

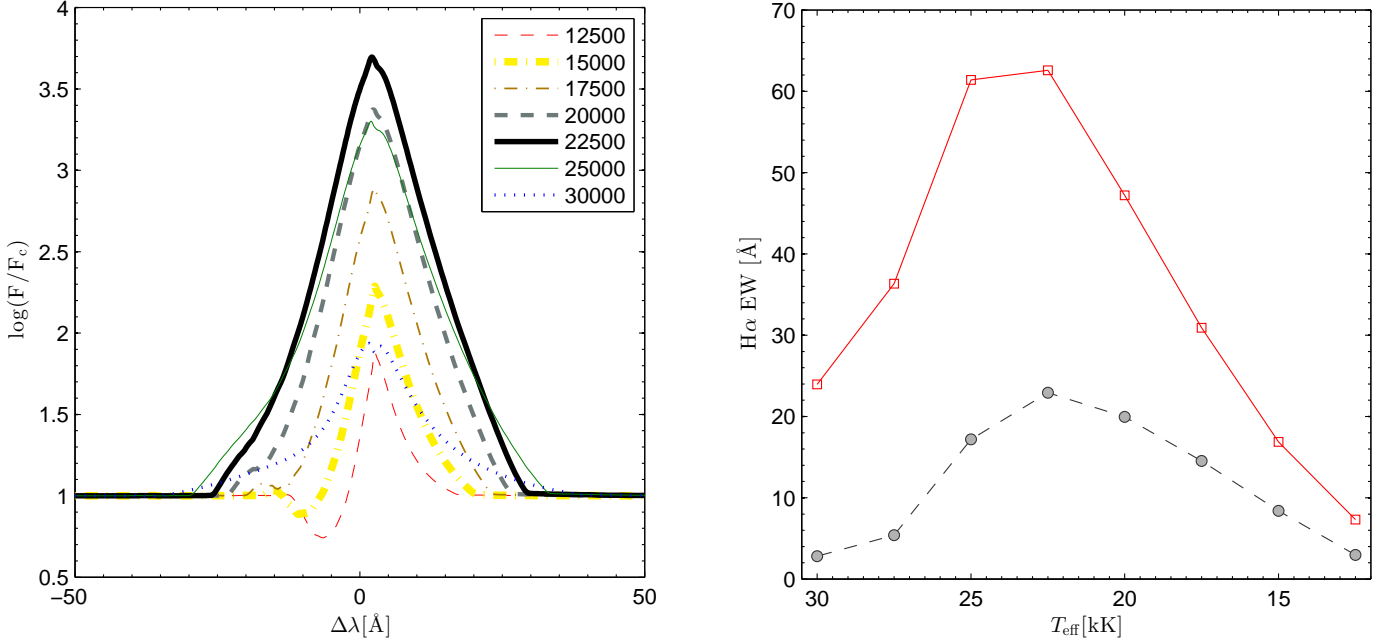


Fig. 14. *Left:* synthetic H α line profiles from clumped models with volume filling factor $f_{\infty} = 0.1$. *Right:* H α line EW as a function of the effective temperature for homogeneous (circles) and clumped (squares) models.

line-forming region is indicated by thick lines. White squares illustrate at which point 50% of the line EW is already formed (see Appendix A). According to the left plot most of the H α photons emerge from regions in which the line is optically thin. Interestingly, when T_{eff} drops from 30 000 to 22 500 K, H α becomes optically thinner. Below 22 500 K, the line changes its behaviour and becomes optically thicker with decreasing T_{eff} .

Figure 15 illustrates that the introduction of micro-clumping would increase the H α optical depth. It seems that at the bi-stability jump ($\sim 22\,500$ K) clumping has the largest impact on the H α optical depth: it increases by an *order of magnitude* at the location where 50% of the line EW is formed (indicated by white squares in Fig. 15). At this point all homogeneous models are optically thin (left panel), while the clumped models are predominantly optically thick in H α (right panel). It appears that the character of H α changes when T_{eff} drops: from an optically thin to an optically thick line (Although, we kept the mass-loss rate and $v_{\infty}/v_{\text{esc}}$ constant in all models). Moreover, it is well known from the observations that $v_{\infty}/v_{\text{esc}}$ drops from 2.6 to 1.3 across the bi-stability jump (Lamers et al. 1995; Markova & Puls 2008). This should produce even sharper differences in the H α line on both sides of the bi-stability jump, as the higher velocity ratio on the hot side is expected to decrease the mean density and thus H α line optical depth, whilst the lower velocity ratio on the cool side would favour higher optical depth of H α line. It is worth mentioning that the inclusion of clumping in the models is equivalent to an increase of the mass-loss rate by factor of $\sqrt{f_{\infty}^{-1}}$ (in comparison to a homogeneous model)³. Because of this, Fig. 15 also illustrates how the H α line optical depth would change if the mass-loss rate were increased by factor of $\sqrt{f_{\infty}^{-1}}$.

All models discussed here have $\dot{M} = 2.33 \times 10^{-6} M_{\odot} \text{yr}^{-1}$ which is nearly 3 times larger than the predicted mass-loss rate for Bsgs around the bi-stability jump. Thus, our homogeneous models are similar to models with roughly 3 times lower mass-

loss rates, but with clumped winds (with $f_{\infty} = 0.1$). Hence, in the context of the predictions, we expect an increase of the H α line optical depth around the bi-stability jump analogous to what is shown in Fig. 15 (i.e., increased by an order of magnitude or more if the mass-loss rate is increased by a factor 5; Vink et al. 1999).

4.3. Impact of macro-clumping on H α

Our finding implies that although changes in the H α optical depth are largest at the cool branch, the H α line EW is less sensitive to micro-clumping in the coolest models. The reasons for that are as follows:

- the H α optical depth increases when T_{eff} is reduced.
- the Ly α optical depth also increases.
- The second level is prevented to recombine to the ground state.

As a consequence, the second level behaves as an effective ground state. Thus, on the cool branch the H α line behaves as an optically thick resonance line and is therefore more sensitive to macro-clumping.

We thus conclude that macro-clumping could play a major role in H α line formation on the cool side of the bi-stability jump. This could significantly affect empirically derived mass-loss rates. Šurlan et al. (2013) showed that for O-type supergiants macro-clumping may resolve the discrepancy between empirical mass-loss rates derived from H α and PV diagnostics (see also Sundqvist et al. 2011).

5. Discussion

In this theoretical study, we have made progress in our understanding of the H α line over the hot and cool branches around the bi-stability jump. What we have not yet discussed is whether our results have a direct bearing on the reported discrepancies of

³ We have tested this assumption by comparing H α optical depth in homogeneous model to a clumped model with factor of $\sqrt{f_{\infty}^{-1}}$ lower \dot{M} .

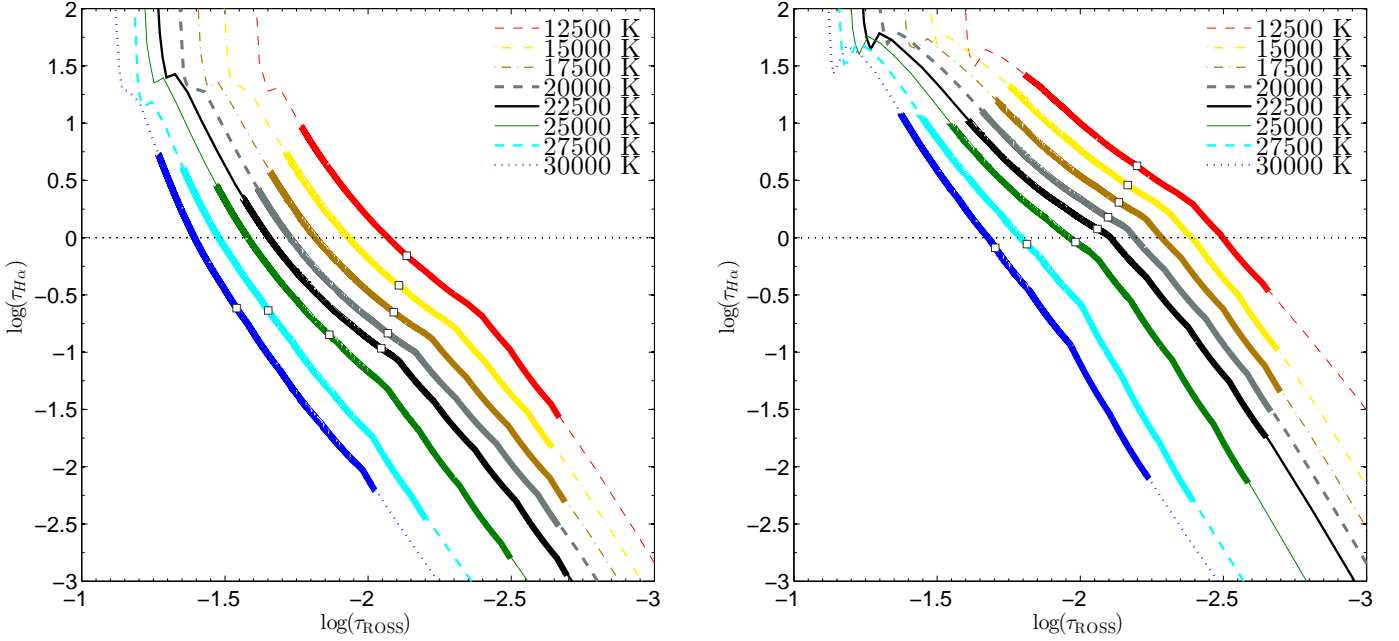


Fig. 15. H α Sobolev optical depth as a function of τ_{ross} for homogeneous (left) and clumped (right) models. Sites where most of the emergent H α photons originates from are set with thick solid lines. White squares represent the point in which 50% of the line EW is already formed (see Appendix A).

empirical late Bsg mass-loss rates from model atmosphere analyses and Monte Carlo iron-line driving calculations of Vink *et al.* (2000). In other words, the problem of the general trend that $\dot{M}_{\text{vink}} > \dot{M}_{\text{H}\alpha}$ for B1 and later supergiants.

Previous investigators, in particular Trundle *et al.* (2004); Trundle & Lennon (2005) and Crowther *et al.* (2006) argued that wind clumping would make the discrepancy between $\dot{M}_{\text{H}\alpha}$ and \dot{M}_{vink} worse, but this conjecture relies on 3 assumptions: (i) that the character of the wind clumps (e.g. optically thin versus optically thick) would remain the same at the bi-stability jump; (ii) that the amount of wind clumping would remain the same; and (iii) that the diagnostic effects of micro-clumping on the H α line would be constant with T_{eff} . How likely is it that the wind structure remains the same when the physics changes at the bi-stability jump?

In the previous sub-section, we noted that our results are in contradiction to those assumptions. The first assumption is also in contradiction to the results of Prinja & Massa (2010), who concluded on the basis of SiIV lines that the Bsgs winds contain optically *thick* clumps, and Oskinova *et al.* (2007) have shown for O star winds that if clumps become optically thick – and the wind becomes porous – clumping might in fact *underestimate* empirical mass-loss rates.

Our results show that H α may become optically thick below the bi-stability jump, whilst conversely it may remain optically thin for hotter objects. This might imply that previous empirical mass-loss rates below the bi-stability jump may have been *underestimated*, whilst those from hotter stars could be correct or slightly overestimated – as a result of micro-clumping. It might also be relevant that the modelled H α lines in the work of Trundle *et al.* (2004) and Trundle & Lennon (2005) do not reproduce the observed H α line shapes. This suggests that the underlying model used in these analyses might not be correct. Moreover, the “derived” values for β are much higher than predicted, and may be an artefact of an inappropriate modelling procedure in case macro-clumping would be relevant. It would

thus be worthwhile for future investigations if the effects seen in Figs. 14 and 15 could indeed explain the reported mass-loss discrepancies.

6. Conclusions

The behaviour of the H α line over T_{eff} range between 30 000 and 12 500 K might be characterised by the management between two processes. Whilst the “rise” is the result of simple recombination (n_3^{\uparrow}), the “fall” is due to the intricate behaviour of the second level (n_2^{\uparrow}). As T_{eff} drops below 22 500 K, the existence of a cool branch may be summarised as follows:

- the high Lyman continuum optical depth makes ionisation from the first level unlikely.
- Ly α becomes optical thick.
- The drain from the second level is suppressed.
- At the coolest model level 2 diverges from higher levels, and it operates like a ground state.
- H α changes its character and behaves like a scattering line with a P-Cygni profile.

During the transition from a recombination to a scattering line, the EW decreases, because recombination lines have a larger (and basically unlimited EW, if the mass-loss rate is increased), whilst a scattering line is confined in its EW as it is dominated by the velocity field. Thus, the EW has to decrease when the line starts to change its character, i.e over the cool branch.

The qualitatively similar H α behaviour including just H, and H+He only models, and metal-blanketed models suggests that the H α behaviour is not related to He or metal properties. Intriguingly, we also found that the effect of clumping on H α is largest at T_{eff} around the bi-stability jump and this is not related to the iron ionisation. Furthermore, if the mass-loss rates are increased at the bi-stability jump (as predicted), then the effect (together with clumping) on H α optical depth would be even stronger. Therefore it is expected that H α changes its character

from an optically thin to an optically thick line and the micro-clumping approximation may no longer be justified.

Although all codes include the physics explained in this work (except for macro-clumping), it is interesting that independent of model complexity, the H α EW peaks at the location of the bi-stability jump for all our models. This might have consequences for both the physics of the bi-stability mechanism, as well as the derived mass-loss rates from H α line profiles, as objects located below the H α EW peak are predicted to be weaker for a similar mass-loss rate, i.e. higher empirical mass-loss rates are required to reproduce a given H α EW if the star is located at a T_{eff} below the peak. Whether this deeper understanding of H α EW over the bi-stability regime would indeed lead to a resolution of the BSG problem remains to be shown with detailed comparisons of our models and observed H α profiles.

This may also be relevant for magnetic field confinement by magnetic spots that could be induced by the sub-surface convection zone in Bsgs (Cantiello *et al.* 2009), as Shultz *et al.* (2013) noted that a pair of spots could have remained undetected by current field searches if the higher theoretical mass-loss rates are employed, but that even the largest spots are ruled out if the lower H α rates are correct.

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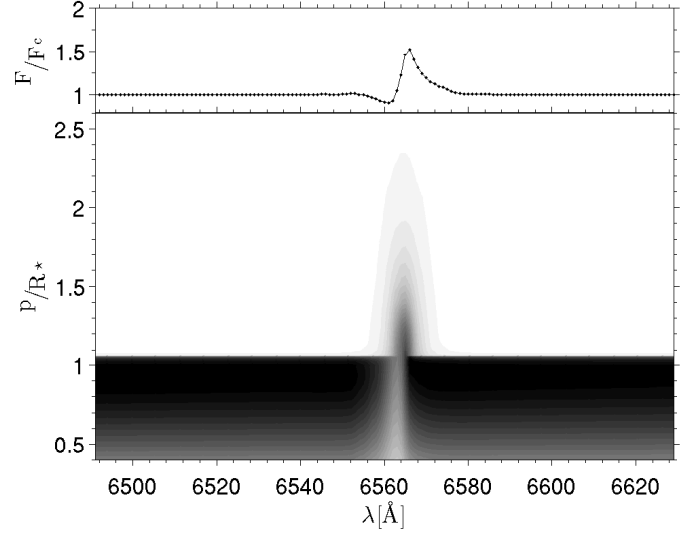


Fig. A.1. Model C ($T_{\text{eff}}=12\,500\text{ K}$). *Bottom*: grey scale plot of the flux like quantity $p \times I(p)$ as a function of impact parameter p/R_* , where R_* is hydrostatic radius. The figure provides the distribution of the emergent intensity around H α from different p . *Top*: corresponding normalised flux in H α , directly obtained by integrating $p \times I(p)$ over the range of p .

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Appendix A: Where in the wind do H α photons originate from?

In order to understand the H α line-formation we show in Fig. A.1 a typical distribution of the emergent intensity $I(p)$, which is scaled by the impact parameter p (see Dessart & Hillier 2005, for details). The top panel of the figure represents the synthetic line profile, which, for each wavelength corresponds to the integral over all p of the scaled intensity $I(p)$ presented at the lower panel. Hence, we are able to identify the contribution at each p to the total line flux.

In other words the figure provides information where in the wind most of the emergent H α photons originate from. This knowledge enables us to display the “evolution” of the H α line (EW) if larger p values are added (or removed) in Fig. A.2. Note that in Fig. A.1 the absorption component emanates for $p/R_* \lesssim 1$ (front of the stellar disc), which is in agreement with the conventional mechanism for P Cygni line formation.

In Fig. A.2 we present how the H α line EW changes when the emergent flux at larger p/R_* (corresponding to $x = r/R_*$ or τ_{ross}) is added. Note that the H α EW is nearly constant for $\log(\tau_{\text{ross}}) > -1.5$: in that range the line is in absorption, mainly produced by the wind in front of the stellar disc ($p/R_* \lesssim 1$), i.e. the H α photons originate at larger distances. From this figure, we

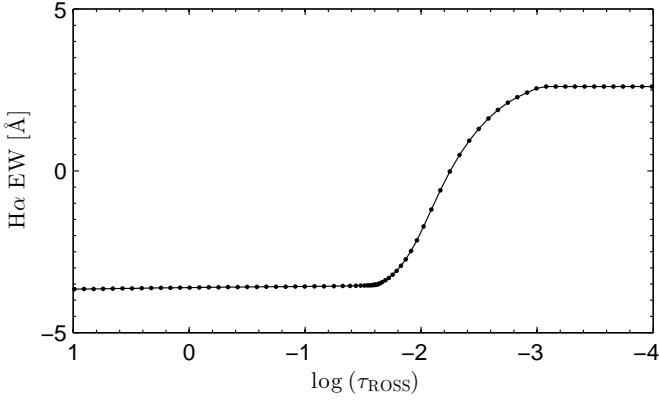


Fig. A.2. H α line EW as a function of τ_{ROSS} . The figure illustrates how the EW changes when outer layers of the star ($p/R_{\star} < 1$) are added.

define the H α line-formation region as the region in which H α changes its EW from 10% to 90%. Although it is by no means conclusive that the line forms in this region, most of the H α photons (in the observer's frame) are emitted from this part of the wind, and the behaviour of the line should depend on the local conditions in that region. Therefore, we investigated the H α related quantities ((n_3/n_2) ratio, $\tau_{\text{Ly}\alpha}$, $\tau_{\text{H}\alpha}$) at this side of the wind in different models.

Appendix B: Model atom (sophisticated models)

Table B.1. For each ion the number of full levels, super levels, and bound-bound transitions are provided.

Ion	Super levels	Full levels	b-b transitions
H I	20	30	435
He I	45	69	905
He II	22	30	435
C I	81	142	3426
C II	40	92	903
C III	51	84	600
C IV	59	64	1446
N I	52	104	855
N II	45	85	556
N III	41	82	578
N IV	44	76	497
N V	41	49	519
O I	32	161	2138
O II	54	123	1375
O III	88	170	1762
O IV	38	78	561
O V	32	56	314
O VI	25	31	203
Si II	9	16	37
Si III	33	33	92
Si IV	22	33	185
P IV	30	90	656
P V	16	62	561
S III	24	44	196
S IV	51	142	1504
S V	31	98	798
Fe I	9	33	47
Fe II	275	827	23004
Fe III	104	1433	57972
Fe IV	74	540	13071
Fe V	50	220	2978
Fe VI	44	433	8662
Fe VII	29	153	1247